

EQUITORSION CONFORM MAPPINGS OF GENERALIZED RIEMANNIAN SPACES

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Abstract. We define an equitorsion conform mapping of two generalized Riemannian spaces and obtain some invariant geometric objects of this mapping, generalizing the tensor of conform curvature.

0. Introduction

A generalized Riemannian space GR_N in the sense of Eisenhart's definition [5] is a differentiable N -dimensional manifold, equipped with nonsymmetric basic tensor g_{ij} .

The use of non-symmetric basic tensor and non-symmetric connection became especially actual after appearance of the works of A. Einstein [1]–[4] related to creation of the Unified Field Theory (UFT). Remark that at UFT the symmetric part $\underline{g_{ij}}$ of the basic tensor g_{ij} is related to the gravitation, and antisymmetric one $\overset{\vee}{g_{ij}}$ to the electromagnetism. M Prvanović [14] and S. Minčić [8] gave geometric interpretations of the torsion and curvature tensors of non-symmetric affine connection.

Consider two N -dimensional generalized Riemannian spaces GR_N and $G\bar{R}_N$. Generalized Cristoffel's symbols of the first kind of the space GR_N and $G\bar{R}_N$ are given by

$$\Gamma_{i,jk} = \frac{1}{2}(g_{ji,k} - g_{jk,i} + g_{ik,j}) \quad \text{and} \quad \bar{\Gamma}_{i,jk} = \frac{1}{2}(\bar{g}_{ji,k} - \bar{g}_{jk,i} + \bar{g}_{ik,j}), \quad (0.1)$$

where, for example, $g_{ij,k} = \partial g_{ij} / \partial x^k$. Connection coefficients of these spaces are generalized Cristoffel's symbols of the second kind $\Gamma_{jk}^i = g^{ip} \Gamma_{p,jk}$ and $\bar{\Gamma}_{jk}^i =$

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$\bar{g}^{jp}\bar{\Gamma}_{p,jk}$, respectively, where $(g^{ij}) = (g_{ij})^{-1}$ and ij denote symmetrisation with division by indices i and j . Generally it is $\Gamma_{jk}^i \neq \Gamma_{kj}^i$. We suppose that $g = \det(g_{ij}) \neq 0$, $\bar{g} = \det(\bar{g}_{ij}) \neq 0$, $\underline{g} = \det(g_{ij}) \neq 0$, $\underline{\bar{g}} = \det(\bar{g}_{ij}) \neq 0$.

One says that a reciprocal one-valued mapping $f : GR_N \rightarrow G\bar{R}_N$ is conform if for the basic tensors g_{ij} and \bar{g}_{ij} of these spaces the condition

$$\bar{g}_{ij} = e^{2\psi} g_{ij} \quad (0.2)$$

is satisfied, where ψ is an arbitrary function of x 's, and the spaces are considered in the common by this mapping system of local coordinates x^i . In this case for the Cristoffel's symbols of the first kind of the spaces GR_N and $G\bar{R}_N$ the relation

$$\bar{\Gamma}_{i,jk} = e^{2\psi} (\Gamma_{i,jk} + g_{ji}\psi_{,k} - g_{jk}\psi_{,i} + g_{ik}\psi_{,j}) \quad (0.3)$$

holds true, and for the Cristoffel's symbols of the second kind

$$\bar{\Gamma}_{jk}^i = \Gamma_{jk}^i + g^{ip}(g_{jp}\psi_{,k} - g_{jk}\psi_{,p} + g_{pk}\psi_{,j}) \quad (0.4)$$

holds. Let us denote $\psi_k = \psi_{,k} = \partial\psi/\partial x^k$ and $\psi^i = g^{ip}\psi_{,p}$. Now from (0.4) we have

$$\bar{\Gamma}_{jk}^i = \Gamma_{jk}^i + g^{ip}(g_{jp}\psi_k - g_{jk}\psi_p + g_{pk}\psi_j) + g^{ip}(g_{jp}\psi_k - g_{jk}\psi_p + g_{pk}\psi_j),$$

i.e.

$$\bar{\Gamma}_{jk}^i = \Gamma_{jk}^i + \delta_j^i \psi_k + \delta_k^i \psi_j - \psi^i g_{jk} + \xi_{jk}^i, \quad (0.5)$$

where

$$\xi_{jk}^i = g^{ip}(g_{jp}\psi_k - g_{jk}\psi_p + g_{pk}\psi_j) = -\xi_{kj}^i \quad (0.5')$$

and $\underset{\vee}{ij}$ denotes an antisymmetrisation with division. In the corresponding points $M(x)$ and $\bar{M}(x)$ of conform mapping we can put

$$\bar{\Gamma}_{jk}^i = \Gamma_{jk}^i + P_{jk}^i \quad (i, j, k = 1, \dots, N), \quad (0.6)$$

where P_{jk}^i is the deformation tensor of the connection Γ of GR_N according to the conform mapping $f : GR_N \rightarrow G\bar{R}_N$.

Notice that in GR_N we have

$$\Gamma_{ip}^p = 0, \quad (0.7)$$

(eq. (2.10) in [13]).

In a generalized Riemannian space one can define four kinds of covariant derivatives [10, 11]. For example, for a tensor a_j^i in GR_N we have

$$\begin{aligned} a_{j|m}^i &= a_{j,m}^i + \Gamma_{pm}^i a_j^p - \Gamma_{jm}^p a_p^i, \\ a_{j|m}^i &= a_{j,m}^i + \Gamma_{mp}^i a_j^p - \Gamma_{mj}^p a_p^i, \\ a_{j|m}^i &= a_{j,m}^i + \Gamma_{pm}^i a_j^p - \Gamma_{mj}^p a_p^i, \\ a_{j|m}^i &= a_{j,m}^i + \Gamma_{mp}^i a_j^p - \Gamma_{jm}^p a_p^i. \end{aligned}$$

Denote by $\left| \cdot \right|_{\theta}$ a covariant derivative of the kind θ in GR_N and $G\bar{R}_N$ respectively. We have [7]

$$g_{ij}|_{\theta} ma = \bar{g}_{ij}|_{\bar{\theta}} ma \equiv 0.$$

In the case of the space GR_N we have five independent curvature tensors [9] (in [9] R_5 is denoted by \tilde{R}_2):

$$\begin{aligned} R_1^i{}_{jmn} &= \Gamma_{jm,n}^i - \Gamma_{jn,m}^i + \Gamma_{jm}^p \Gamma_{pn}^i - \Gamma_{jn}^p \Gamma_{pm}^i, \\ R_2^i{}_{jmn} &= \Gamma_{mj,n}^i - \Gamma_{nj,m}^i + \Gamma_{mj}^p \Gamma_{np}^i - \Gamma_{nj}^p \Gamma_{mp}^i, \\ R_3^i{}_{jmn} &= \Gamma_{jm,n}^i - \Gamma_{nj,m}^i + \Gamma_{jm}^p \Gamma_{np}^i - \Gamma_{nj}^p \Gamma_{pm}^i + \Gamma_{nm}^p (\Gamma_{pj}^i - \Gamma_{jp}^i), \\ R_4^i{}_{jmn} &= \Gamma_{jm,n}^i - \Gamma_{nj,m}^i + \Gamma_{jm}^p \Gamma_{np}^i - \Gamma_{nj}^p \Gamma_{pm}^i + \Gamma_{mn}^p (\Gamma_{pj}^i - \Gamma_{jp}^i), \\ R_5^i{}_{jmn} &= \frac{1}{2} (\Gamma_{jm,n}^i + \Gamma_{mj,n}^i - \Gamma_{jn,m}^i - \Gamma_{nj,m}^i + \Gamma_{jm}^p \Gamma_{pn}^i + \Gamma_{mj}^p \Gamma_{np}^i \\ &\quad - \Gamma_{jn}^p \Gamma_{mp}^i - \Gamma_{nj}^p \Gamma_{pm}^i). \end{aligned}$$

We use the conform mapping $f: GR_N \rightarrow G\bar{R}_N$ to obtain tensors $\bar{R}_{\theta}^i{}_{jmn}$ ($\theta = 1, \dots, 5$), where for example

$$\bar{R}_1^i{}_{jmn} = \bar{\Gamma}_{jm,n}^i - \bar{\Gamma}_{jn,m}^i + \bar{\Gamma}_{jm}^p \bar{\Gamma}_{pn}^i - \bar{\Gamma}_{jn}^p \bar{\Gamma}_{pm}^i. \quad (0.8)$$

In the case of conform mapping $f: R_N \rightarrow \bar{R}_N$ of Riemannian spaces R_N and \bar{R}_N [6, 15] we have an invariant geometric object

$$C^i{}_{jmn} = R^i{}_{jmn} + \delta_m^i P_{jn} - \delta_n^i P_{jm} + P_m^i g_{jn} - P_n^i g_{mj} \quad (0.9)$$

where

$$P_{jm} \equiv \frac{1}{N-2} (R_{jm} - \frac{1}{2(N-1)} R g_{jm}),$$

and $R^i{}_{jmn}$ is Riemann-Cristoffel's curvature tensor of the space R_N , R_{jm} Ricci's tensor and R a scalar curvature.

The object $C^i{}_{jmn}$ is called a conform curvature tensor [6, 15]. Having a conform mapping of two generalized Riemannian spaces, we cannot find a generalization of the tensor of conform curvature as an invariant of conform mapping in general case. For that reason we define a special conform mapping.

A mapping $f: GR_N \rightarrow G\bar{R}_N$ is an equitorsion conform mapping if the torsion tensors of the spaces GR_N and $G\bar{R}_N$ are equal. Then from (0.5) and (0.6) we have

$$\xi_{jk}^i = 0. \quad (0.10)$$

In [12] we have investigated equitorsion geodesic mappings of generalized Riemannian spaces.

1. Equitorsion conform curvature tensor of the first kind

Using (0.6), we get a relation between the first kind curvature tensors of the spaces GR_N and $G\bar{R}_N$ [12, 16]

$$\bar{R}_{1jmn}^i = R_{1jmn}^i + P_{jm|n}^i - P_{jn|m}^i + P_{jm}^p P_{pn}^i - P_{jn}^p P_{pm}^i + 2\Gamma_{mn}^p P_{jp}^i.$$

Substituting P with respect to (0.5, 6, 10), and using (0.7'), we obtain

$$\begin{aligned} \bar{R}_{1jmn}^i &= R_{1jmn}^i + \delta_j^i (\psi_{m|n} - \psi_{n|m}) + \delta_m^i (\psi_{j|n} - \psi_j \psi_n) \\ &\quad - \delta_n^i (\psi_{j|m} - \psi_j \psi_m) - (\psi_{j|n}^i - \psi_n \psi_j^i) g_{jm} + (\psi_{j|m}^i - \psi_m \psi_j^i) g_{jn} \\ &\quad - \delta_n^i \psi^p \psi_p g_{jm} + \delta_m^i \psi^p \psi_p g_{jn} + 2\delta_j^i \Gamma_{mn}^p \psi_p + 2\Gamma_{mn}^i \psi_j - 2\Gamma_{j.mn} \psi^i. \end{aligned} \quad (1.1)$$

Denoting

$$\psi_{ij} = \psi_{i|j} - \psi_i \psi_j, \quad \psi_j^i = g^{ip} \psi_{pj} \quad (1.2a)$$

$$\Delta_1 \psi = g^{pq} \psi_p \psi_q = \psi_p \psi^p \quad (1.2b)$$

and using the relation

$$\psi_{mn} - \psi_{nm} = -2\Gamma_{mn}^p \psi_p \quad (1.3)$$

in (1.1), we get

$$\begin{aligned} \bar{R}_{1jmn}^i &= R_{1jmn}^i + \delta_m^i \psi_{jn} - \delta_n^i \psi_{jm} + \psi_m^i g_{jn} - \psi_n^i g_{jm} \\ &\quad + (\delta_m^i g_{jn} - \delta_n^i g_{jm}) \Delta_1 \psi + 2\Gamma_{mn}^i \psi_j - 2\Gamma_{j.mn} \psi^i. \end{aligned} \quad (1.4)$$

Further, let us denote

$$\Delta_2 \psi = g^{pq} \psi_p |q \quad (1.5)$$

Then we have

$$\psi_p^p = \psi_{pq} g^{pq} = (\psi_{p|q} - \psi_p \psi_q) g^{pq} = \Delta_2 \psi - \Delta_1 \psi.$$

Contracting by indices i and n in (1.4) we get

$$\bar{R}_{1jm} = R_{1jm} - (N-2) \psi_{jm} - [\Delta_2 \psi + (N-2) \Delta_1 \psi] g_{jm} - 2\Gamma_{j.mp} \psi^p. \quad (1.6)$$

From (0.2) we get

$$\bar{g}^{ij} = e^{-2\psi} g^{ij}. \quad (1.7)$$

In (1.6) multiplying by g^{jm} and contracting by j and then by m we get

$$e^{2\psi} \bar{R} = R - 2(N-1) \Delta_2 \psi - (N-1)(N-2) \Delta_1 \psi, \quad (1.8)$$

where $\bar{R}_1 = \bar{g}^{pq}\bar{R}_{pq}$, and $R_1 = g^{pq}R_{pq}$ are scalar curvature of the first kind of the spaces $G\bar{R}_N$ and GR_N respectively. From (1.8) we have

$$\Delta_2\psi = \frac{1}{2(N-1)}(R_1 - e^{2\psi}\bar{R}_1) - \frac{N-2}{2}\Delta_1\psi. \quad (1.9)$$

Substituting (1.9) in (1.6) we get

$$\begin{aligned} (N-2)\psi_{1jm} = & R_{1jm} - \bar{R}_{1jm} - \frac{1}{2(N-1)}(R_1 - e^{2\psi}\bar{R}_1)g_{jm} \\ & - \frac{N-2}{2}\Delta_1\psi g_{jm} - 2\Gamma_{j.m\check{p}}\psi^{\check{p}}. \end{aligned} \quad (1.10)$$

Let us denote in the space GR_N

$$P_{1jm} \equiv \frac{1}{N-2}(R_{1jm} - \frac{1}{2(N-1)}Rg_{jm}) \quad (1.10')$$

and analogously \bar{P}_{1jm} in the space $G\bar{R}_N$. In this case for ψ_{jm} we obtain

$$\psi_{1jm} = P_{1jm} - \bar{P}_{1jm} - \frac{1}{2}\Delta_1\psi g_{jm} - \frac{2}{N-2}\Gamma_{j.m\check{p}}\psi^{\check{p}}. \quad (1.11)$$

Substituting (1.11) in (1.4), we get

$$\begin{aligned} \bar{R}_{1jmn}^i = & R_{1jmn}^i + \delta_m^i(P_{1jn} - \bar{P}_{1jn}) - \delta_n^i(P_{1jm} - \bar{P}_{1jm}) \\ & + P_{1m}^i g_{jn} - \bar{P}_{1m}^i \bar{g}_{jn} - P_{1n}^i g_{jm} + \bar{P}_{1n}^i \bar{g}_{jm} \\ & + \frac{2}{N-2}(\delta_n^i \Gamma_{j.m\check{p}} - \delta_m^i \Gamma_{j.n\check{p}} + \Gamma_{n\check{p}}^i g_{jm} - \Gamma_{m\check{p}}^i g_{jn})\psi^{\check{p}} \\ & + 2\Gamma_{mn}^i \psi_j - 2\Gamma_{j.m\check{p}}\psi^{\check{p}}. \end{aligned} \quad (1.12)$$

We can see that it follows from (0.2)

$$\psi_i = \frac{1}{2N}(\frac{\partial}{\partial x^i} \ln \bar{g} - \frac{\partial}{\partial x^i} \ln g) \quad (1.13)$$

where $g = \det(g_{ij})$, $\bar{g} = \det(\bar{g}_{ij})$. From (0.10) and (1.13) we obtain

$$\Gamma_{j.nm}^i \psi^i = \frac{1}{2N}\bar{\Gamma}_{j.nm}^i \bar{g}^{ip} \frac{\partial}{\partial x^p} \ln \bar{g} - \frac{1}{2N}\Gamma_{j.nm}^i g^{ip} \frac{\partial}{\partial x^p} \ln g \quad (1.14)$$

and

$$\Gamma_{qn}^i g_{mj} \psi^q = \frac{1}{2N}\bar{\Gamma}_{qn}^i \bar{g}_{mj} \bar{g}^{pq} \frac{\partial}{\partial x^p} \ln \bar{g} - \frac{1}{2N}\Gamma_{qn}^i g_{mj} g^{pq} \frac{\partial}{\partial x^p} \ln g. \quad (1.15)$$

Taking into account (1.13, 14,15), we can write the relation (1.12) in the form

$$\bar{C}_1^i{}_{jmn} = C_1^i{}_{jmn}, \quad (1.16)$$

where

$$\begin{aligned}
C_1^i{}_{jmn} &= R_1^i{}_{jmn} + \delta_m^i P_{jn} - \delta_n^i P_{jm} + P_{1m}^i g_{jn} - P_{1n}^i g_{jm} \\
&+ \frac{1}{N(N-2)} (\delta_m^i \Gamma_{j.np} - \delta_n^i \Gamma_{j.mp} + \Gamma_{mp}^i g_{jn} - \Gamma_{np}^i g_{jm}) g^{pq} \frac{\partial}{\partial x^q} \ln g \\
&+ \frac{1}{N} (\Gamma_{j.mn} g^{ip} - \Gamma_{mn}^i \delta_j^p) \frac{\partial}{\partial x^p} \ln g
\end{aligned} \tag{1.17}$$

and analogously for $\overline{C}_1^i{}_{jmn}$. From (1.16) we see that the tensor $C_1^i{}_{jmn}$ is an invariant of equitorsion conform mapping, and one can call it the equitorsion conform curvature tensor of the first kind. So, we have

THEOREM 1. *Let generalized Riemannian spaces GR_N and $G\overline{R}_N$ be defined by virtue of their nonsymmetric basic tensors g_{ij} and \overline{g}_{ij} respectively. The equitorsion conform curvature tensor of the first kind $C_1^i{}_{jmn}$ (1.17) is an invariant of the equitorsion conform mapping $f: GR_N \rightarrow G\overline{R}_N$, defined by (0.2), (0.5), (0.10), i.e. (1.16) is in force, where the tensor P_1 is given by (1.10').*

2. Equitorsion conform curvature tensor of the second kind

For the second kind curvature tensors of the spaces GR_N and $G\overline{R}_N$ we get the relation [12, 16]

$$\overline{R}_2^i{}_{jmn} = R_2^i{}_{jmn} + P_{m_j}^i |_{n_2} - P_{n_j}^i |_{m_2} + P_{m_j}^p P_{np}^i - P_{n_j}^p P_{mp}^i + 2\Gamma_{nm}^p P_{pj}^i,$$

i.e., using (0.5,6,10) one obtains

$$\begin{aligned}
\overline{R}_2^i{}_{jmn} &= R_2^i{}_{jmn} + \delta_m^i \psi_{jn} - \delta_n^i \psi_{jm} + \psi_{2m}^i g_{nj} - \psi_{2n}^i g_{mj} \\
&+ (\delta_m^i g_{mj} - \delta_n^i g_{mj}) \Delta_1 \psi + 2\Gamma_{nm}^i \psi_j - 2\Gamma_{nm}^p \psi^i g_{pj},
\end{aligned} \tag{2.1}$$

where

$$\psi_{ij} = \psi_{i|j} - \psi_i \psi_j, \quad \psi_j^i = g^{ip} \psi_{pj}, \quad \Delta_1 \psi = g^{pq} \psi_p \psi_q \tag{2.2}$$

Now, analogously to previous case, we get the invariant object of the equitorsion conform mapping $f: GR_N \rightarrow G\overline{R}_N$

$$\begin{aligned}
C_2^i{}_{jmn} &= R_2^i{}_{jmn} + \delta_m^i P_{jn} - \delta_n^i P_{jm} + P_{2m}^i g_{jn} - P_{2n}^i g_{mj} \\
&+ \frac{1}{N(N-2)} (\delta_m^i \Gamma_{j.pn} - \delta_n^i \Gamma_{j.pm} + \Gamma_{pm}^i g_{jn} - \Gamma_{pn}^i g_{mj}) g^{pq} \frac{\partial}{\partial x^q} \ln g \\
&+ \frac{1}{N} (\Gamma_{j.nm} g^{ip} - \Gamma_{nm}^i \delta_j^p) \frac{\partial}{\partial x^p} \ln g
\end{aligned} \tag{2.3}$$

where

$$P_{2jm} \equiv \frac{1}{N-2} (R_{2jm} - \frac{1}{2(N-1)} R_2 g_{mj}), \quad (2.4)$$

R_{2jm} is Ricci's curvature tensor of the second kind and R_2 is a scalar curvature tensor of the second kind. The object C_{2jmn}^i is a tensor and we call it equitortion conform curvature tensor of the second kind. Accordingly, we have

THEOREM 2. *Starting from the curvature tensor R_{2jmn}^i , under conditions as in Theorem 1, one obtains an invariant tensor C_{2jmn}^i (2.3) of the equitortion conform mapping of generalized Riemannian spaces, where P_2 is given according to (2.4).*

3. Equitortion conform curvature tensor of the third kind

In the case of the third kind curvature tensors of the spaces GR_N and $G\bar{R}_N$ we get the relation [12, 16]

$$\begin{aligned} \bar{R}_{3jmn}^i &= R_{3jmn}^i + P_{jm|n}^i - P_{nj|m}^i + P_{jm}^p P_{np}^i - P_{nj}^p P_{pm}^i \\ &\quad + 2P_{nm}^p \Gamma_{pj}^i + 2P_{nm}^p \Gamma_{pj}^i \end{aligned}$$

i.e., because of (0.5,6,10), (1.2a,b) and (2.2),

$$\begin{aligned} \bar{R}_{3jmn}^i &= R_{3jmn}^i + \delta_m^i \psi_{jn} - \delta_n^i \psi_{jm} - \psi_{2n}^i g_{jm} + \psi_{1m}^i g_{nj} \\ &\quad + (\delta_m^i g_{nj} - \delta_n^i g_{jm}) \Delta_1 \psi + 2\psi_m \Gamma_{nj}^i + 2\psi_n \Gamma_{mj}^i - 2\psi^p g_{nm} \Gamma_{pj}^i \end{aligned} \quad (3.1)$$

Also, the following is satisfied

$$\psi_{mn} = \psi_{mn} + 2\Gamma_{mn}^p \psi_p, \quad \psi_n^i = \psi_n^i + 2g_{pn}^i \Gamma_{pn}^q \psi_q. \quad (3.2)$$

From (3.1), (3.2) and (0.10) we get

$$\begin{aligned} \bar{R}_{3jmn}^i &= R_{3jmn}^i + \delta_m^i \psi_{jn} - \delta_n^i \psi_{jm} + \psi_{1m}^i g_{nj} - \psi_{1n}^i g_{jm} \\ &\quad + (\delta_m^i g_{nj} - \delta_n^i g_{jm}) \Delta_1 \psi + 2\psi_m \Gamma_{nj}^i + 2\psi_n \Gamma_{mj}^i - 2\psi^p g_{nm} \Gamma_{pj}^i \\ &\quad + 2\delta_m^i \Gamma_{jn}^p \psi_p - 2g_{pn}^i \Gamma_{pn}^q \psi_q g_{jm}. \end{aligned} \quad (3.3)$$

Contracting (3.3) with respect to i and n , and using (1.5), we get

$$\bar{R}_{3jm} = R_{3jm} - (N-2)\psi_{jm} - [\Delta_2 \psi + (N-2)\Delta_1 \psi] g_{jm} - \psi^p \Gamma_{m.pj}. \quad (3.4)$$

Multiplying (3.4) by $\bar{g}^{jm} = e^{-2\psi} g^{jm}$ and contracting we get

$$\Delta_1 \psi = \frac{1}{2(N-1)} (R_3 - e^{2\psi} \bar{R}_3) - \frac{N-2}{2} \Delta_1 \psi. \quad (3.5)$$

Substituting (3.5) in (3.4) and denoting

$$P_{\underset{3}{j}m} = \frac{1}{N-2}(\underset{3}{R}_{jm} - \frac{1}{2(N-1)}\underset{3}{R}g_{jm}) \quad (3.6)$$

in GR_N and analogously in $G\overline{R}_N$, in this case for ψ_{jm} we obtain

$$\psi_{\underset{1}{j}m} = P_{\underset{3}{j}m} - \overline{P}_{\underset{3}{j}m} - \frac{1}{2}\Delta_1\psi g_{jm} - \frac{2}{N-2}\Gamma_{m,pj}\psi^p. \quad (3.7)$$

Substituting (3.7) in (3.3) and using (1.14,15) we get

$$\overline{C}_{\underset{3}{j}mn}^i = C_{\underset{3}{j}mn}^i \quad (3.8)$$

where

$$\begin{aligned} C_{\underset{3}{j}mn}^i &= R_{\underset{3}{j}mn}^i + \delta_m^i P_{\underset{3}{j}n} - \delta_n^i P_{\underset{3}{j}m} + P_{\underset{3}{m}n\underset{j}{\underline{}}}^i - P_{\underset{3}{n}m\underset{j}{\underline{}}}^i \\ &+ \frac{1}{N(N-2)}(\delta_m^i \Gamma_{n,pj} - \delta_n^i \Gamma_{m,pj})g^{pq} \frac{\partial}{\partial x^q} \ln g \\ &+ \frac{1}{N}(g^{ip} \Gamma_{pn}^q g_{jm} - \delta_m^q \Gamma_{nj}^i - \delta_n^q \Gamma_{mj}^i \\ &+ \Gamma_{pj}^i g_{nm} g^{pq} - \delta_m^i \Gamma_{jn}^q) \frac{\partial}{\partial x^q} \ln g \end{aligned} \quad (3.9)$$

And analogously for $\overline{C}_{\underset{3}{j}mn}^i$ of the space $G\overline{R}_N$. From (3.8) we can see that the tensor $C_{\underset{3}{j}mn}^i$ is an invariant of equitorsion conform mapping, and one can call it the equitorsion conform curvature tensor of the third kind. Now we have

THEOREM 3. *From the curvature tensor $R_{\underset{3}{j}mn}^i$, under the conditions as in Theorem 1, we obtain an invariant tensor $C_{\underset{3}{j}mn}^i$ (3.9) of the equitorsion conform mapping $f: GR_N \rightarrow G\overline{R}_N$, where $P_{\underset{3}{j}}$ is given according to (3.6).*

4. Equitorsion conform curvature tensor of the fourth kind

For curvature tensors of the fourth kind we get [12, 16]

$$\begin{aligned} \overline{R}_{\underset{4}{j}mn}^i &= R_{\underset{4}{j}mn}^i + P_{jm|n}^i - P_{nj|m}^i + P_{jm}^p P_{np}^i - P_{nj}^p P_{pm}^i \\ &+ 2P_{mn}^p \Gamma_{pj}^i + 2P_{mn}^p P_{pj}^i \end{aligned}$$

i.e.

$$\begin{aligned} \overline{R}_{\underset{4}{j}mn}^i &= R_{\underset{4}{j}mn}^i + \delta_m^i \psi_{jn} - \delta_n^i \psi_{jm} + \psi_{\underset{1}{m}n\underset{j}{\underline{}}}^i - \psi_{\underset{1}{n}m\underset{j}{\underline{}}}^i \\ &+ (\delta_m^i g_{nj} - \delta_n^i g_{jm})\Delta_1\psi + 2\psi_n \Gamma_{mj}^i + 2\psi_m \Gamma_{nj}^i - 2\psi^p g_{mn} \Gamma_{pj}^i \\ &+ 2\delta_m^i \Gamma_{jn}^p \psi_p - 2g^{ip} \Gamma_{pn}^q \psi_q g_{jm}. \end{aligned}$$

In this case, analogously to previous case, we get an invariant object of the equitorsion conform mapping in the form

$$C_4^i{}_{jmn} = R_4^i{}_{jmn} + \delta_m^i P_4{}_{jn} - \delta_n^i P_4{}_{jm} + P_4^i{}_{m\underline{gnj}} - P_4^i{}_{n\underline{gjm}} + \frac{1}{N(N-2)}(\delta_m^i \Gamma_{n\underline{pj}} - \delta_n^i \Gamma_{m\underline{pj}})g^{pq} \frac{\partial}{\partial x^q} \ln g \quad (4.1)$$

$$+ \frac{1}{N}(g^{ip} \Gamma_{pn\underline{gjm}} - \delta_m^q \Gamma_{nj} - \delta_n^q \Gamma_{mj} + \Gamma_{pj\underline{gnm}} g^{pq} - \delta_m^i \Gamma_{jn}^q) \frac{\partial}{\partial x^q} \ln g,$$

$$P_4{}_{jm} = \frac{1}{N-2}(R_4{}_{jm} - \frac{1}{2(N-1)}R_4 g_{jm}), \quad (4.2)$$

where $R_4{}_{jm}$ is Ricci's curvature tensor of the fourth kind and R_4 a scalar curvature of the fourth kind. The object $C_4^i{}_{jmn}$ is a tensor and we call it equitorsion conform curvature tensor of the fourth kind of the equitorsion conform mapping. So, the next theorem is valid.

THEOREM 4. *From the curvature tensor $R_4^i{}_{jmn}$, under the conditions as in Theorem 1, one obtains an invariant tensor $C_4^i{}_{jmn}$ (4.1) of the equitorsion conform mapping of generalized Riemannian spaces, where P is given with respect to (4.2).*

5. Equitorsion conform curvature tensor of the fifth kind

For the curvature tensors of the fifth kind of the spaces GR_N and $G\overline{R}_N$ we find the relation [12, 16]

$$\overline{R}_5^i{}_{jmn} = R_5^i{}_{jmn} + \frac{1}{2}(P_{jm|n}^i - P_{jn|m}^i + P_{mj|n}^i - P_{nj|m}^i + P_{jm}^p P_{pn}^i - P_{jn}^p P_{mp}^i + P_{mj}^p P_{np}^i - P_{nj}^p P_{pm}^i)$$

i.e.

$$\begin{aligned} \overline{R}_5^i{}_{jmn} = & R_5^i{}_{jmn} + \frac{1}{2}[\delta_m^i (\psi_{j|n} + \psi_{j|n} - 2\psi_j \psi_n) - \delta_n^i (\psi_{j|m} + \psi_{j|m} - 2\psi_j \psi_m) \\ & + (\psi_{|m}^i + \psi_{|m}^i - 2\psi_m \psi^i)g_{jn} - (\psi_{|n}^i + \psi_{|n}^i - 2\psi_n \psi^i)g_{mj} \\ & + 2(\delta_m^i g_{jn} - \delta_n^i g_{jm})\psi_p \psi^p]. \end{aligned} \quad (5.1)$$

Let us denote

$$\psi_{34}{}_{jn} = \frac{1}{2}(\psi_{j|n} + \psi_{j|n} - 2\psi_j \psi_n), \quad \psi_{34}^i{}_j = g^{ip} \psi_{pj}, \quad \Delta_1 \psi = g^{pq} \psi_p \psi_q. \quad (5.2)$$

Then

$$\begin{aligned} \overline{R}_5^i{}_{jmn} = & R_5^i{}_{jmn} + \delta_m^i \psi_{34}{}_{jn} - \delta_n^i \psi_{34}{}_{jm} + \psi_{34}^i{}_m g_{jn} - \psi_{34}^i{}_n g_{mj} \\ & + (\delta_m^i g_{jn} - \delta_n^i g_{jm})\Delta_1 \psi. \end{aligned} \quad (5.3)$$

Contracting by indices i, n and denoting

$$\overline{R}_{5jmp}^p = \overline{R}_{5jm}, \quad R_{5jmp}^p = R_{5jm}, \quad \Delta_{34}\psi = \frac{1}{2}g_{53}^{pq}(\psi_{p|q} + \psi_{p|q}), \quad (5.4)$$

we obtain

$$\overline{R}_{5jm} = R_{5jm} - (N-2)\psi_{34}^j - [\Delta_{34}\psi + (N-2)\Delta_1\psi]g_{jm}, \quad (5.5)$$

wherefrom, multiplying by $\overline{g}^{jm} = e^{-2\psi}g^{jm}$ and contracting by j and then by m one obtains

$$\Delta_{34}\psi = \frac{1}{2(N-1)}(R_{55} - e^{2\psi}\overline{R}_{55}) - \frac{N-2}{2}\Delta_1\psi. \quad (5.6)$$

From (5.5) and (5.6) we get

$$\psi_{34}^j = P_{5jm} - \overline{P}_{5jm} - \frac{1}{2}\Delta_1\psi g_{jm} \quad (5.7)$$

where we denoted

$$P_{5jm} = \frac{1}{N-2}(R_{5jm} - \frac{1}{2(N-1)}R_{55}g_{jm}) \quad (5.8)$$

in GR_N and analogously \overline{P}_{5jm} in $G\overline{R}_N$.

Analogously to previous cases eliminating ψ_{34}^j from (5.3) we can write

$$\overline{C}_{5jmn}^i = C_{5jmn}^i, \quad (5.9)$$

where we denoted

$$C_{5jmn}^i = R_{5jmn}^i + \delta_m^i P_{5jn} - \delta_n^i P_{5jm} + P_{5m}^i g_{nj} - P_{5n}^i g_{jm}. \quad (5.10)$$

The object C_{5jmn}^i is an invariant of the equitorsion conform mapping. We call it equitorsion conform curvature tensor of the fifth kind. So, we have

THEOREM 5. *Starting from the curvature tensor R_{5jmn}^i , under the conditions as in the Theorem 1, we obtain an invariant tensor C_{5jmn}^i (5.10) of the equitorsion mapping $f: GR_N \rightarrow G\overline{R}_N$, where P is given according to (5.8).*

If $GR_N(G\overline{R}_N)$ reduces to $R_N(\overline{R}_N)$, then the objects $C_{\theta jmn}^i$ ($\theta = 1, \dots, 5$) reduce to the conform curvature tensor (0.9).

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